

General Relativistic Simulations of Cosmic Large Scale Structure

based on arXiv:1308.6524

and work in progress with C. Clarkson, E. DiDio, R. Durrer, and M. Kunz

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W3: Origin of Cosmic Structures

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Motivation: A Case for GR

- Newton vs. Einstein
- The Issue of Backreaction
- Relativistic Sources of Stress-Energy

The Framework

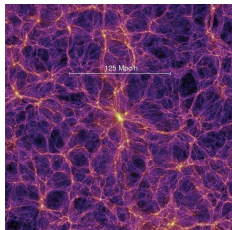
- Choice of Variables
- Weak Field Approximation
- System of Equations
- Algorithmic Solutions

Numerical Results

- A Plane-Symmetric Setup

A Case for GR

Newton vs. Einstein



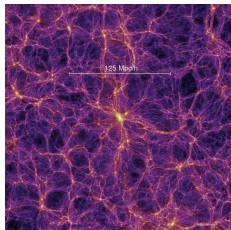
Millennium simulation,
Springel *et al.* 2005

In order to study the regime of nonlinear structure formation, large N-body simulations are the method of choice.

N-body simulations use
Newton's law of gravity

A Case for GR

Newton vs. Einstein



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N-body simulations use
Newton's law of gravity

Works well for *nonrelativistic matter* (CDM), because of
see, e.g., Green & Wald 2012

- Exact correspondence between Newtonian gravity and GR on the background solution (FRW)
- Exact correspondence also on the level of linear (scalar) perturbations
- Nonlinear scale \ll Hubble scale

A Case for GR

Newton vs. Einstein (cont.)

The Newtonian picture has several drawbacks, though

- Strong assumption about material content of the Universe
- Misses some degrees of freedom (gravity waves!)
- Gauge issues are not apparent
- Trivial propagation of light beams (relativistic effects have to be put back “by hand”)

A Case for GR

Newton vs. Einstein (cont.)

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- Strong assumption about material content of the Universe
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A unified relativistic treatment of structure formation would automatically solve these issues. When constructing observables (galaxy catalogs, lensing maps etc.), all geometric effects and gauge issues would be treated in a transparent way.

A Case for GR

The Issue of Backreaction

Long standing question: how important is nonlinear evolution of structure for understanding & interpreting the observed “average” cosmological evolution?

References include Ellis 1984, Buchert 2000 & 2008, Schwarz 2002 & 2012, Wetterich 2003, Kolb, Matarrese, Notari & Riotto 2005, Buchert & Ellis 2005, Räsänen 2011

A Case for GR

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This issue has many facets. Some can be addressed in the Newtonian picture, others require a relativistic treatment (→ perturbation theory, exact solutions . . .).

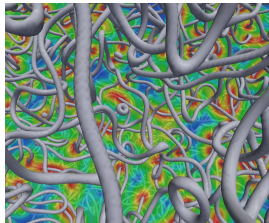
A unified relativistic treatment of structure formation would be the logical framework to address the issue in full generality.

A Case for GR

Relativistic Sources of Stress-Energy

GR effects are expected to be important for intrinsically relativistic entities

- Cosmic strings
- Dynamical Dark Energy
- Relativistic particles (neutrinos?)
- ...



credit: *Daverio et al.*

In order to test some of the proposed alternatives/extensions to Λ CDM, general relativistic simulations may be necessary in order to obtain percent accuracy required by future observations (e.g. Euclid)

The Framework

Choice of Variables

Metric of perturbed FRW in “longitudinal gauge”

$$ds^2 = -(1 + 2\Psi) dt^2 + a^2(t)[(1 - 2\Phi) \delta_{ij} + h_{ij}] dx^i dx^j - 2B_i dx^i dt$$

Gauge condition: $\nabla^i B_i = \nabla^i h_{ij} = \delta^{ij} h_{ij} = 0$

The Framework

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Stress-energy tensor

$$T_\nu^\mu = \bar{T}_\nu^\mu + \delta T_\nu^\mu, \quad \bar{T}_\nu^\mu = \text{diag}(-\bar{\rho}, \bar{P}, \bar{P}, \bar{P}), \quad \bar{P} = w\bar{\rho}$$

Fix “background” equation of state for each constituent \rightarrow
background scale factor a solves Friedmann’s equations for \bar{T}_ν^μ

The Framework

Weak Field Approximation

Perturbative approach:

- Metric perturbations Ψ, Φ, \dots remain small in cosmological context ($\sim 10^{-5}$) \rightarrow keep only to first order
- Spatial derivatives $\Psi_{,i}, \dots$ are $\sim v$ ($\sim 10^{-3}$) \rightarrow keep to quadratic order
- Second spatial derivatives $\Delta\Psi, \dots$ are $\sim \delta$ and therefore non-perturbative

Green & Wald 2012, J.A., Daverio, Durrer & Kunz 2013

The Framework

System of Equations

“ $G_0^0 = 8\pi GT_0^0$ ”:

$$\frac{1}{a^2} (1 + 4\Phi) \Delta\Phi - 3H\dot{\Phi} - 3H^2\Psi + \frac{3}{2a^2} (\nabla\Phi)^2 = -4\pi G\delta T_0^0$$

The Framework

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$$“G_i^i - 3G_0^0 - \frac{1}{H}\dot{G}_0^0 = 8\pi G(T_i^i - 3T_0^0 - \frac{1}{H}\dot{T}_0^0)”:$$

$$\begin{aligned} & (1 + 2\Phi - 2\Psi) \Delta\Psi - (\nabla\Psi)^2 - \nabla\Psi\nabla\Phi + \\ & \frac{1}{H}\partial_t \left[\Delta\Phi + 4\Phi\Delta\Phi + \frac{3}{2} (\nabla\Phi)^2 \right] = \\ & 4\pi G\frac{a}{H} \left[\delta T_{0,i}^i - \delta T_0^i \left(3\Phi_{,i} - \Psi_{,i} + a\dot{B}_i \right) \right. \\ & \left. - a\dot{\Phi} \left(3\delta T_0^0 - \delta T_i^i \right) - \frac{a}{2} \delta^{ik} \dot{h}_{jk} \delta T_i^j \right] \end{aligned}$$

The Framework

System of Equations (cont.)

“ $G_i^0 = 8\pi GT_i^0$ ”:

$$-\frac{4}{a^2}\Delta B_i - \frac{1}{a}\dot{\Phi}_{,i} - \frac{H}{a}\Psi_{,i} = 4\pi G\delta T_i^0$$

The Framework

System of Equations (cont.)

$$“G_i^0 = 8\pi G T_i^0”:$$

$$-\frac{4}{a^2}\Delta B_i - \frac{1}{a}\dot{\Phi}_{,i} - \frac{H}{a}\Psi_{,i} = 4\pi G\delta T_i^0$$

$$“G_j^i - \frac{1}{3}\delta_j^i G_k^k = 8\pi G(T_j^i - \frac{1}{3}\delta_j^i T_k^k)”:$$

$$\begin{aligned} & \ddot{h}_{ij} + 3H\dot{h}_{ij} - \frac{1}{a^2}\Delta h_{ij} + \frac{1}{a}\left(\dot{B}_{(i,j)} + 2HB_{(i,j)}\right) \\ & + \frac{1}{a} \text{“traceless part } [(1 + 4\Phi)\Phi_{,ij} - (1 + 2\Phi - 2\Psi)\Psi_{,ij} \\ & + \Psi_{,i}\Psi_{,j} - 2\Phi_{(i}\Psi_{,j)} + 3\Phi_{,i}\Phi_{,j}]” = 8\pi G\left(\delta_{ik}\delta T_j^k - \frac{1}{3}\delta_{ij}\delta T_k^k\right) \end{aligned}$$

The Framework

System of Equations (cont.)

In order to close the system of equations, one needs evolution equations for all sources of stress-energy.

The Framework

System of Equations (cont.)

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Geodesic equation for nonrelativistic massive particles

$$\dot{v}^i + H v^i + \delta^{ij} \left(\frac{1}{a} \Psi_{,j} - \dot{B}_j - H B_j + \frac{2}{a} B_{[j,k]} v^k \right) = 0$$

determines the evolution of the particle ensemble and therefore the evolution of the full T_{ν}^{μ} of CDM.

The Framework

Algorithmic Solutions

Φ : parabolic equation (diffusion type)

- Explicit scheme too inefficient (Courant condition!)
- First-order (in time) implicit scheme shows excellent performance in 1D tests
- ADI (*Alternating Direction Implicit*) scheme in 3D should perform well (easy to implement & parallelizable)

The Framework

Algorithmic Solutions

Φ : parabolic equation (diffusion type)

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Ψ : elliptic equation

- Nonlinear Gauß-Seidel / Multigrid solver shows excellent performance in 1D tests
- Same class of solvers is already used for the Poisson equation in modern Newtonian codes

The Framework

Algorithmic Solutions (cont.)

B_i : linear elliptic operator

Two possibilities:

- Solve in Fourier space (transverse component can easily be extracted, but incompatible with AMR)
- Use Multigrid solver (gauge condition more difficult to implement)

The Framework

Algorithmic Solutions (cont.)

B_i : linear elliptic operator

Two possibilities:

- Solve in Fourier space (transverse component can easily be extracted, but incompatible with AMR)
- Use Multigrid solver (gauge condition more difficult to implement)

h_{ij} : linear hyperbolic equation (wave equation)

- No conceptual problem, but can be expensive (depending on relevant range of scales)
- h_{ij} does not enter the geodesic equation for massive particles (at our approximation order) \rightarrow expendable, leave for future work

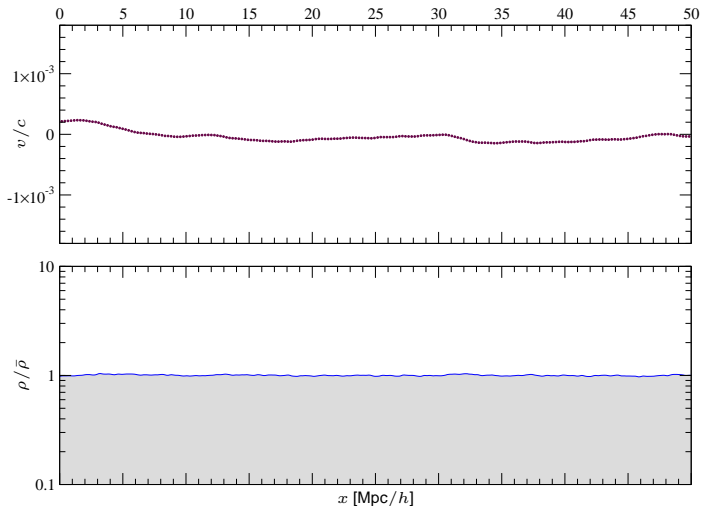
Numerical Results

A Plane-Symmetric Setup

- Restriction to plane-symmetric configuration ($y - z$ -plane) trivializes two dimensions \rightarrow high resolution possible with cheap computational requirements (no parallelization)
- No vector & tensor perturbations (by construction)
- 32768 particles, 4096 grid points
- Initial conditions: Gaussian random field obtained from semi-realistic initial power spectrum
- Initialized at $z > 1000$ using linear theory (Zel'dovich approximation)

Numerical Results

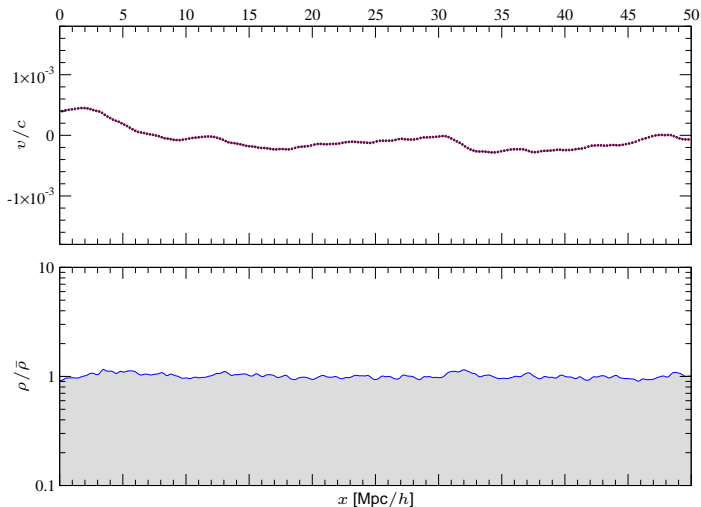
A Plane-Symmetric Setup



$z = 100$

Numerical Results

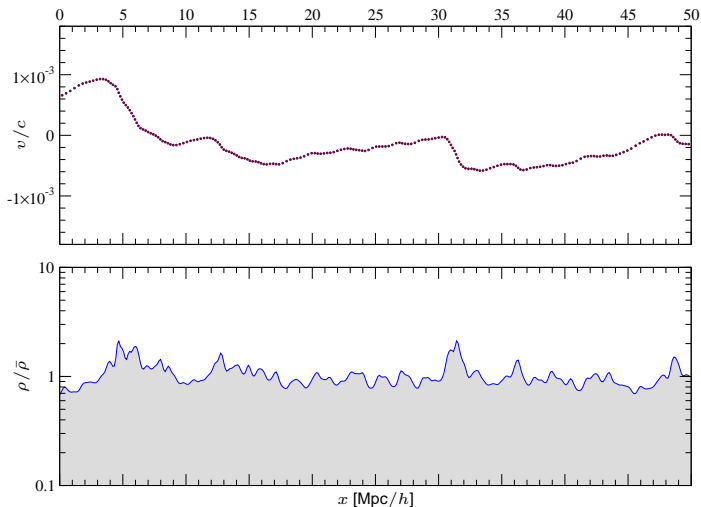
A Plane-Symmetric Setup



$z = 26$

Numerical Results

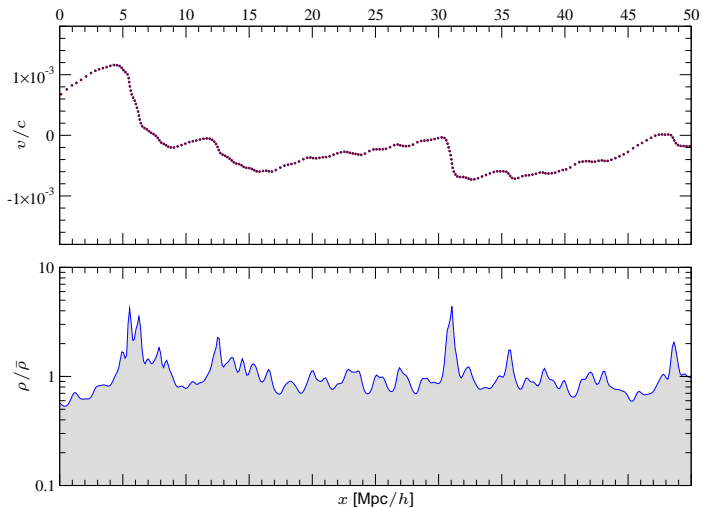
A Plane-Symmetric Setup



$z = 5.3$

Numerical Results

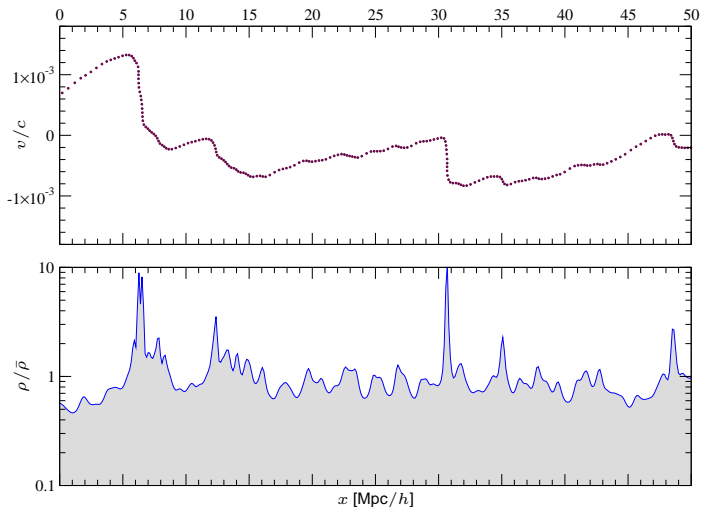
A Plane-Symmetric Setup



$z = 3.0$

Numerical Results

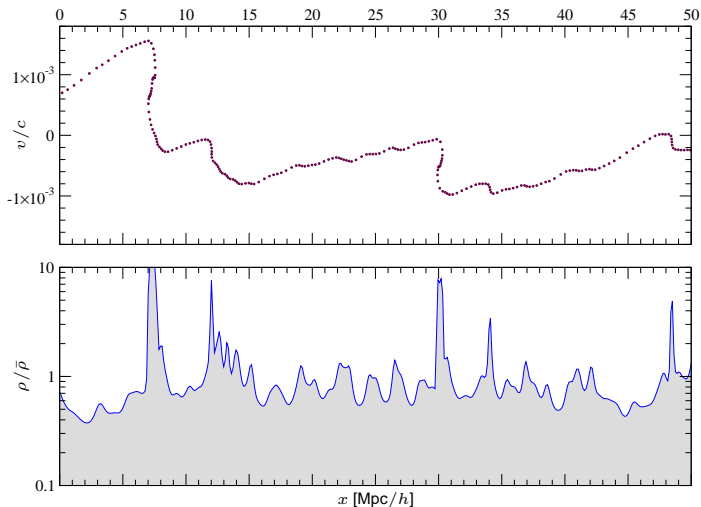
A Plane-Symmetric Setup



$z = 2.0$

Numerical Results

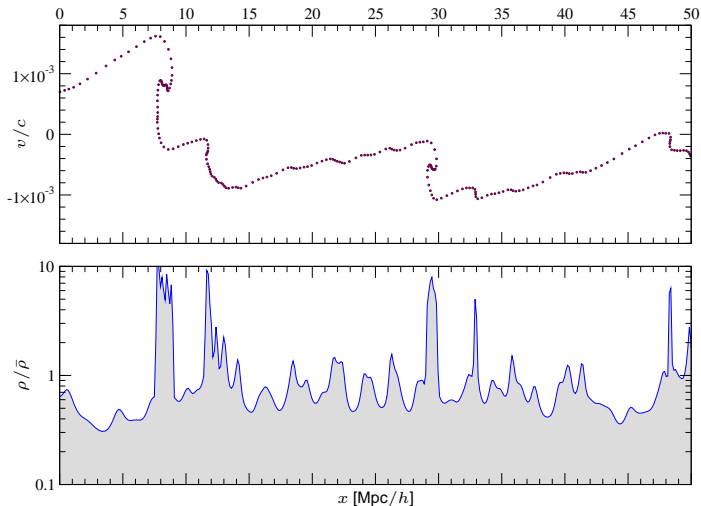
A Plane-Symmetric Setup



$z = 1.0$

Numerical Results

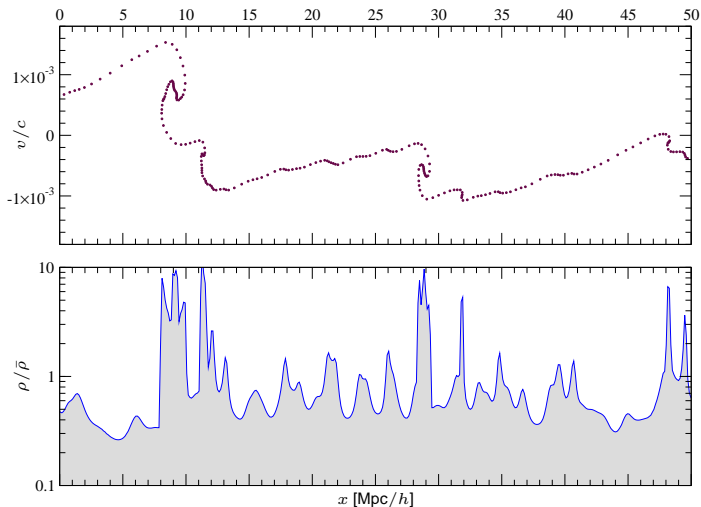
A Plane-Symmetric Setup



$z = 0.4$

Numerical Results

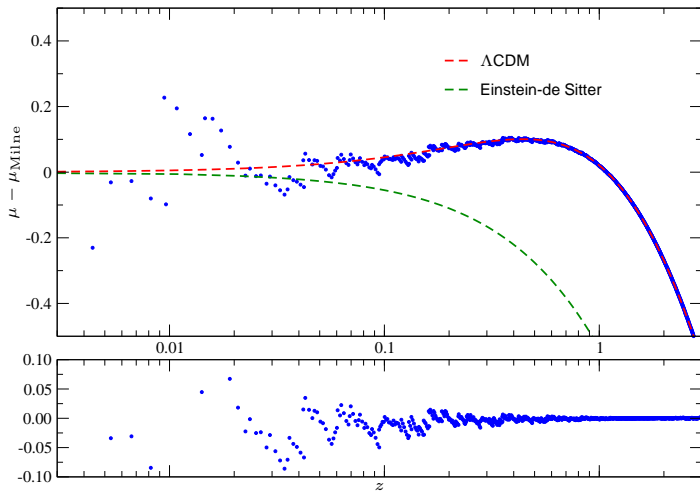
A Plane-Symmetric Setup



$z = 0.0$

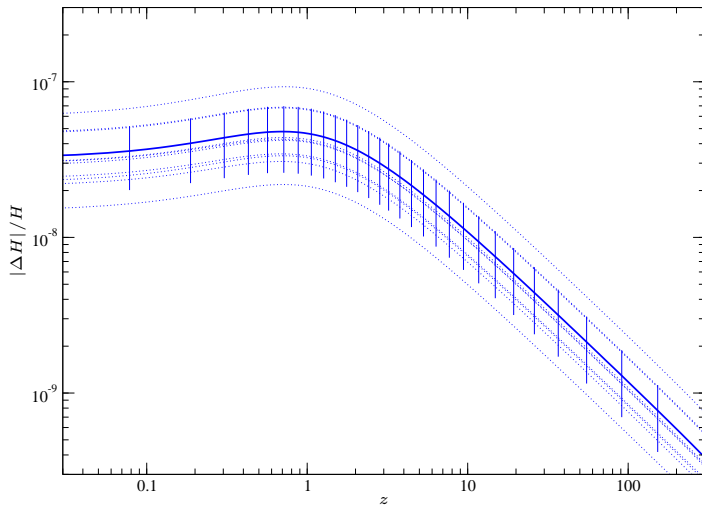
Numerical Results

Luminosity Distance



Numerical Results

Newtonian vs. GR Simulation



Fin

Summary

- Cosmological simulations within a **GR framework** are feasible
- A unified relativistic treatment is a clear, logical and transparent way to address the **most general observables** with minimal assumptions about the cosmological model
- Technology is useful for simulations with **relativistic sources** (dynamical DE, cosmic strings, neutrinos) – feasibility depends on ability to model the sources accurately
- For CDM simulations, modifications are computationally relatively inexpensive (but may be unnecessary)
- The issue of **backreaction** can be addressed quantitatively within the non-perturbative regime